

VAGAS, Istvan

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kozlonyn 41 no.6:458 D'61

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1. "Hidrologiai Kozlony" szerkeszto bizottsagi tagja es  
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ny. tanar.; VAGAS, Istvan,okl. mernok.

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1. Epitoipari es Kozlekedesi Muszaki Egyetem (for Szilagyi).
2. A Vizgazdalkodasi Tudomanyos Kutato Intezet tud.munkatarsa  
(for Vagas).

VAGAS, Istvan

HUNGARY

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Tudomanyos Kutato Intezet), Budapest

Budapest, Hidrologiai Kozlony, No 5, Oct 62, pp 399-410.

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Kezleny" szerkeszto bizottsagi tagja.

VAGAS, Istvan; LASZLOFFY, Woldemar, dr.; IVICSIKS, Lajos, dr.; UBELL, Karoly,  
dr.; OLLÓS, Geza

Society and technical news. Hidrologiai Kozlony 42 no.1:18,29,36,71,75,  
79, 80 F '62.

1. "Hidrologiai Kozlony" szerkeszto bizottsagi tagja (for Vagas and  
Laszlóffy). 2. "Hidrologiai Kozlony" feszerkesztoje (for Ollós).

VAGAS, Istvan

"Handbook of hydraulic engineering structures" by Odon Staro-  
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204,216,232,266 Jl '62.

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Reviewed by Istvan Vagas. Hidrologiai kozlony 42 no.3:293  
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VAGAS, Istvan

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381,389,393,398,449,452 O '62.

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VAGAS, Istvan; OLLOS, Geza

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470, 477, 487 D '62.

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(for Ollos).

HANKO, Zoltan; VAGAS, Istvan; PERENYI, Karoly; MUSZKALAY, Laszlo

Society and technical news. Hidrelegiai kozlony 43 no.1:6, 18, 30,  
38, 76, 81 F '63.

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VAGAS, Istvan; SZABO, Laszlo, dr.; KONTUR, Gyorgy; IVICSICS, Lajos, dr.;  
OLLOS, Geza

Society and technical news. Hidrologiai kozlony 43 no.2:98,112,118,  
121, 129, 164, 184 Ap '63.

1. "Hidrologiai Kozlony" szerkeszto bizottsagi tagja (for Ivicsics and  
Vagas). 2. "Hidrolegiai Kozlony" feszerkesztoje (for Ollos).

VAGAS, Istvan; KONTUR, Gyorgy; IVICSIOS, Ferenc; BAUER, Jeno; LOVAS,  
Laszlo

Society and technical news. Hidrologiai kozlony 43 no.3:197, 204,  
218, 224, 228, 236, 250, 265 Je '63.

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of water level swelling. 318-322

I. Vizgazdalkodasi Tudomanyos Kutato Intezet, Budapest; "Hidro-  
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VAGAS, Istvan; AUJESZKY, Geza

Society and technical news. Hidrologiai kozlony 43 no.4:307,  
317 Ag'63.

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VAGAS, Istvan; KOZAK, Imre

Society and technical news. Hidrologia kozlony 43 no.6:520,  
527 D '63.

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rologiai Kozlony" szerkeszto bizottsagi tagja es rovatvezetoje  
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48 Ja'64.

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means of an Euclidean model. Ibid.:206-208 My '64.

1. Scientific Research Institute of Water Resources Develop-  
ment; editorial board member, "Hidrologiai Kozlony."

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Je '64.

1. Editorial board member, "Hidrologiai Kozlony" (for Vagas).

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395-419 S '64.

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VAGAS, Istvan; OLLOS, Geza

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1. Chair of Water Resources Development at the Technical University  
of Building and Transportation, Budapest (for Bozoky-Szezich).
2. Editorial Board Member, "Mikrologiai Kozlony" (for Vagas).

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VAGAUNESCU, G.; SIMU, C.

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1. Clinica I medicala Cluj si Catedra de anatomie patologica Cluj.

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ZAGREANU, I., dr.; VAGAUNESCU, Gh., dr.; SUCIU, I., dr.

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14 no.4:621-624 My '62.

1. Lucrare efectuata in Clinica I medicala, Cluj, director acad. A. Moga.  
(CARDIOVASCULAR DISEASES) (RURAL HEALTH)  
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VAGAUNESCU, Gh. dr.

Influence of iodine treatment on some humoral changes in patients with atherosclerosis. Med. intern. (Bucur.) 10 no.5: 591-600 My'64

1. Lucrare efectuata in Clinica I medicala I.M.F. [ Institutul medico-farmaceutic], Cluj (director: acad. A.Moga).

VAGAYTSEV, V.I., tekhnik po defektoskopii (g.Chita)

Skillful operators of defectoscopes. Put' i put.khoz. 6  
no.11:35 '62. (MIRA 16:1)  
(Railroads—Maintenance and repair)

VAGAYTSEV, V.I., tekhnik po defektoskopii (Chita)

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(MIRA 16:1)  
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USSR- / Virology. Human and Animal Viruses

E-2

Abs Jour: Ref Zhur - Biol., No 6, 1958, 23955

Author : Vagazhanova, V. A.

Inst : Not given

Title : Interaction of Tumors and Viruses.

Orig Pub: Vopr. virusologii, 1957, No 3, 179-181

**Abstract:** When a mixture of Erlich ascitic carcinoma with viruses of tick encephalitis, encephalomyelitis, encephalomyocarditis, herpes, chick-pest and fixated rabies virus are administered to mice intracerebrally, the percentage of mouse deaths after a two-hour contact was higher than on administration of these viruses and ascitic carcinoma separately. In intraperitoneal administration of the enumerated viruses, as well as influenza and ectromelia viruses, the percentage of mouse

Card 1/2

USSR / Virology. Human and Animal Viruses

E-2

Abs Jour: Ref Zhur - Biol., No 6, 1958, 23955

Abstract: deaths after a two hour contact differed but  
little from the control.

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L 43995-66 EWT(1) GW  
ACC NR: AT6021504

(N)

SOURCE CODE: UR/2531/86/000/187/0003/0012

AUTHOR: Vager, B. G.; Zilitinkevich, S. S.

ORG: none \*

21  
B71

TITLE: Method of calculating the height of the lower boundary of clouds based on numerical forecasting data

SOURCE: Leningrad. Glavnaya geofizicheskaya observatoriya. Trudy, no. 187, 1966.  
Fizika pogranichnogo sloya atmosfery (Physics of the atmospheric boundary layer), 3-12

TOPIC TAGS: atmospheric cloud, cloud cover, cloud level, WEATHER FORECASTING

ABSTRACT: An attempt is made to establish a functional relationship between the height of the lower boundary of the cloud cover and parameters the values of which lend themselves to numerical forecasting. In the mathematical description of the method, the input equations, boundary conditions, and integration of the input equations and additional simplifications which are performed on an electronic computer are cited. Several specific examples, the starting material for which were experimental data obtained from an investigation of the lower cloud cover during the fall of 1962 in the region of Dnepropetrovsk were examined in order to compare the calculated values of the height of the lower boundary of clouds with the observed

Card 1/2

L 43995-66

ACC NR: AT6021504

values. Their mean square deviation amounted to 149 m. Orig. art. has: 2 tables and 39 formulas.

SUB CODE: 04/ SUBM DATE: none/ ORIG REF: 016/ OTH REF: 002

Card 2/2 ULR

L 08153-67 EWT(1) GW

ACC NR: AT6021506

(N)

SOURCE CODE: UR/2531/66/000/187/0044/0053

27  
25  
B+1AUTHOR: Vager, B. G.; Tseytin, G. Kh.ORG: noneTITLE: Structure of the atmospheric boundary layer under stationary conditions  
(nonlinear problem)SOURCE: Leningrad. Glavnaya geofizicheskaya observatoriya. Trudy, no. 187, 1966.  
Fizika pogranichnogo sloya atmosfery (Physics of the atmospheric boundary layer), 44-53

TOPIC TAGS: atmospheric boundary layer, atmospheric turbulence, turbulent diffusion

ABSTRACT: A scheme is presented for computing some meteorological  
structure of the atmospheric boundary layer under stationary  
conditions (nonlinear problem). IN: Glavnaya geofizi-  
cheskaya observatoriya. Fizika pogranichnogo sloya atmosfery  
(Physics of the surface boundary layer of the atmosphere),  
1966, 44-53. (ITS: Trudy, no. 187, 1966).

A scheme is presented for computing some meteorological characteristics of the boundary layer of the atmosphere for various states of thermal stability, with turbulent energy diffusion taken into account. Since diffusion of turbulent energy can be neglected in the special case of neutral stratification, this method can be reduced to a numerical scheme for computation.

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ACC NR: AT6021506

The problem reduces to solving the following system of equations: equations of motion under horizontally homogeneous conditions:

$$\left. \begin{aligned} \frac{d}{dz} h(z) \frac{du}{dz} + 2\omega_s v(z) &= 0 \\ \frac{d}{dz} h(z) \frac{dv}{dz} - 2\omega_s [u(z) - V_g] &= 0 \end{aligned} \right\} \quad (1)$$

the equation of the influx of heat

$$ak(z) \left[ \frac{dT}{dx} + \gamma_p \right] = - \frac{P_0}{c_p \rho} - \frac{R(z) - R(0)}{c_p \rho} \quad (2)$$

and the equations of turbulent energy balance

$$h(z) \left[ \left( \frac{du}{dz} \right)^2 + \left( \frac{dv}{dz} \right)^2 \right] - f ak(z) \left[ \frac{dT}{dx} + \gamma_p \right] - c_0 \frac{b^2(z)}{h(z)} + a_1 \frac{d}{dz} h(z) \frac{db(z)}{ds} = 0 \quad (3)$$

$$h(z) = l(z) \sqrt{b(z)} \quad (4)$$

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The expression for the scale of turbulence  $\lambda(z)$  (according to D. L. Laykhtman and S. S. Zilitinkevich) is

$$\lambda(z) = -2c_0^{\frac{1}{4}} \times \frac{F(z)}{dF(z)}, \quad (5)$$

$$F(z) = \left(\frac{du}{dz}\right)^2 + \left(\frac{dv}{dz}\right)^2 - \frac{ag}{T} \left(\frac{dT}{dz} + \gamma_p\right). \quad (6)$$

The following notation is adopted in equations (1)-(6):  $u(z)$ ,  $v(z)$  are the horizontal components of the wind velocity (the  $x$ -axis is directed along the isobar);  $k(z)$  is the turbulence factor;  $v_g$  is the velocity of the geostrophic wind;  $T(z)$  is the absolute air temperature;  $\gamma_p$  is the equilibrium temperature gradient;  $P_0$  is the value of the turbulent heat flux at ground level;  $R(z)$  is the radiant heat flux;  $b(z)$  is the energy of turbulent pulsations;  $a$  is the ratio of the turbulence factor for heat to the turbulence factor for momentum; and  $c_0$  and  $a_1$  are constants. The remaining notation is standard. In this formulation of the problem, it is considered that the values of the radiation for  $R(z)$  and turbulent flux near the ground  $P_0$  are given. Thus, the system is closed with six equations (1-5) available for six unknowns:  $u(z)$ ,  $v(z)$ ,  $k(z)$ ,  $\lambda(z)$ ,  $b(z)$ , and  $T(z)$ .

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The boundary conditions are as follows:

$$u(z) = v(z) = 0 \text{ when } z = z_0 \quad (7)$$

$$u(z) \rightarrow V_g, \quad v(z) \rightarrow 0, \quad z \rightarrow \infty \quad (8)$$

$$k(z)|_{z=z_0} = x v_* z_0, \quad (9)$$

where  $V_g$  is the dynamic velocity, and

$$v^2 = k(z) \sqrt{\left(\frac{du}{dz}\right)^2 + \left(\frac{dv}{dz}\right)^2} \Big|_{z=z_0} \quad (10)$$

$$k(z) \frac{db(z)}{dz} \Big|_{z=z_0} = 0, \quad (11)$$

$$b(z) \rightarrow 0, \quad z \rightarrow \infty \quad (12)$$

$z_0$  is the surface roughness parameter. The boundary condition (11) means that turbulent energy does not penetrate the underlying surface.

A scheme for solving the problem is given for the general case. First, the dimensionless height and the desired functions are introduced by the formulas:

$$\eta = V_g \int_{z_0}^z \frac{dz}{k(z)} = \frac{2w_*}{V_g} \int_{z_0}^z \frac{dz}{\beta(z)}, \quad (13)$$

$$\beta(z) = \frac{2w_*}{V_g^2} k(z), \quad (14)$$

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ACC NR: AT6021506

$$\left. \begin{array}{l} u_0(\eta) = \frac{u(\eta)}{V_s} \\ v_0(\eta) = \frac{v(\eta)}{V_s} \end{array} \right\}$$

(16)

$$\Phi(\eta) = \frac{c_0^2}{V_s^2} b(\eta).$$

(16)

Then, instead of equations (1), (3), and the boundary conditions (7), (8), (11), and (12), we get

$$\frac{d^2\sigma(\eta)}{d\eta^2} - \beta(\eta)\sigma(\eta) = 0, \quad (17)$$

$$\frac{sd^2\Phi(\eta)}{d\eta^2} - \Phi^2(\eta) + E(\eta) = 0, \quad (18)$$

$$\sigma(\eta)|_{\eta=0} = 1,$$

(19)

$$\sigma(\eta) \rightarrow 0, \quad \eta \rightarrow \infty$$

(20)

$$\Phi(\eta) \rightarrow 0, \quad \eta \rightarrow \infty$$

(21)

$$\left. \frac{d\Phi(\eta)}{d\eta} \right|_{\eta=0} = 0, \quad (22)$$

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Where

$$\bar{O}(\eta) = 1 - u_0(\eta) - t v_0(\eta), \quad (23)$$

$$E(\eta) = \left( \frac{du_0}{d\eta} \right)^2 + \left( \frac{dv_0}{d\eta} \right)^2 - s_1 \beta^2(\eta) \left[ \frac{dT}{dz} + \gamma_p \right] =$$

$$= \left( \frac{du_0}{d\eta} \right)^2 + \left( \frac{dv_0}{d\eta} \right)^2 + s_2 \beta(\eta) P(z), \quad (24)$$

$$P(z) = -\alpha c_p \rho k(z) \left[ \frac{dT}{dz} + \gamma_p \right], \quad (25)$$

$$s = \frac{a_1}{\sqrt{c_0}}, \quad (26)$$

$$s_1 = \frac{\alpha g}{4 \omega_x^2 T}, \quad (27)$$

$$s_2 = \frac{g}{2 \omega_x c_p \rho T V^2}, \quad (28)$$

On the basis of (4), (5), (9), and (10), we find that

$$\beta(\eta) = m_0 \sqrt{E(\eta)} e^{\int_{\eta_0}^{\eta} \gamma_p(z) dz}. \quad (29)$$

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where

$$m_0 = \frac{2u_s x_0}{v_e \sqrt{1 + \frac{x_0}{L}}}, \quad m_0 \approx \frac{2u_s x_0}{v_e} \quad (30)$$

for  $x_0 \ll L$ , where  $L = \frac{g}{T} \cdot \frac{P_0}{c_p P}$  is the height of the

surface boundary sublayer, according to Obukhov and Monin. If stratification is neutral, when  $L = \infty$ , (30) becomes exact. When  $dT/dz$  or  $P(z)$  is given for the case of an atmosphere which is not in a state of equilibrium, the system of equations, in dimensionless form, contains only the two parameters  $m_0$  and  $s_1$  (or  $s_2$ ), and universal constants.

When the state of the atmosphere is neutral ( $P(z) \equiv 0$ ), the system depends on only one parameter  $m_0$  which can be expressed by the Rossby parameter

$$Ro = \frac{V_r}{2u_s x_0} \quad (31)$$

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Thus, the scheme for solving the system can be written, in the general case, as follows:

- 1) An approximate profile of the turbulence factor  $\beta(n)$  is given;
- 2) The equation of motion (17) is solved by the selected  $\beta(n)$  with boundary conditions (19)–(20);
- 3) The profile  $E(n)$  is calculated by formula (24) with  $s_1$  (or  $s_2$ ) given and the profile  $dT/dz$  or  $P(z)$ ;
- 4) The profile  $\phi(n)$  is determined by formula (18), making use of  $E(n)$  and boundary conditions (21)–(22);
- 5) A new profile  $\beta(n)$  is determined by formula (29), using the given profile  $m_0$  and profiles  $\phi(n)$  and  $E(n)$ .

If the new profile differs noticeably from the preceding one, this procedure is repeated with the new profile, starting with step 2).

A simplified numerical method of computation is given for the case of neutral stratification (without taking turbulent-energy diffusion into consideration). In this case, equation (18) takes the form

$$E(n) - \Phi^2(n) = 0, \quad (32)$$

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where

$$E(\eta) = \left( \frac{du_0}{d\eta} \right)^2 + \left( \frac{dv_0}{d\eta} \right)^2 \quad (33)$$

Thus, equation (29) is replaced by

$$\phi(\eta) = m_0 \sqrt{E(\eta)} e^{\int \frac{1}{E(\eta)} d\eta} \quad (34)$$

Reducing the linear boundary problem to Cauchy problems, one can seek the desired function  $\phi(\eta)$  in the form

$$Q(\eta) = Q_1(\eta) - \frac{Q_1(\eta_H)}{Q_2(\eta_H)} Q_2(\eta), \quad (35)$$

where  $Q_1(\eta)$  and  $Q_2(\eta)$  are solutions of equation (17) with the boundary conditions

$$Q_1(\eta)|_{\eta=0} = 1, \quad \frac{dQ_1(\eta)}{d\eta}|_{\eta=0} = 0, \quad (36)$$

$$Q_2(\eta)|_{\eta=0} = 0, \quad \frac{dQ_2(\eta)}{d\eta}|_{\eta=0} = 1, \quad (37)$$

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and  $\eta_0$  is a sufficiently great height. Then

$$\left. \begin{array}{l} O_1(\eta) = A(\eta) + iB(\eta) \\ O_2(\eta) = C(\eta) + iD(\eta) \end{array} \right\} \quad (38)$$

Then, we get instead of (17), (36), and (37):

$$\left. \begin{array}{l} \frac{d^2A(\eta)}{d\eta^2} = -\beta(\eta)B(\eta) \\ \frac{d^2B(\eta)}{d\eta^2} = \beta(\eta)A(\eta) \\ \frac{d^2C(\eta)}{d\eta^2} = -\beta(\eta)D(\eta) \\ \frac{d^2D(\eta)}{d\eta^2} = \beta(\eta)C(\eta) \end{array} \right\} \quad (39)$$

$$\left. \begin{array}{l} A(\eta)|_{\eta=0} = 1, \quad B(\eta)|_{\eta=0} = 0 \\ \frac{dA(\eta)}{d\eta}|_{\eta=0} = 0, \quad \frac{dB(\eta)}{d\eta}|_{\eta=0} = 0 \\ C(\eta)|_{\eta=0} = 0, \quad D(\eta)|_{\eta=0} = 0 \\ \frac{dC(\eta)}{d\eta}|_{\eta=0} = 1, \quad \frac{dD(\eta)}{d\eta}|_{\eta=0} = 0 \end{array} \right\} \quad (40)$$

The wind velocity components  $u_0(\eta)$  and  $v_0(\eta)$  are expressed

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by the formulas

$$\left. \begin{array}{l} u_0(\eta) = 1 - A(\eta) + L(\eta_{II})C(\eta) - \bar{L}(\eta_{II})D(\eta) \\ v_0(\eta) = -B(\eta) + \dot{L}(\eta_{II})D(\eta) + \bar{L}(\eta_{II})C(\eta) \end{array} \right\} \quad (41)$$

where

$$\left. \begin{array}{l} L(\eta_{II}) = \frac{A(\eta_{II})C(\eta_{II}) + B(\eta_{II})D(\eta_{II})}{C(\eta_{II})^2 + D(\eta_{II})^2} \\ \bar{L}(\eta_{II}) = \frac{B(\eta_{II})C(\eta_{II}) - A(\eta_{II})D(\eta_{II})}{C(\eta_{II})^2 + D(\eta_{II})^2} \end{array} \right\} \quad (42)$$

The Adams method was used for numerical integration of the system (39)-(40) with modified coefficients in which differences up to the sixth order inclusive were taken into account. This permitted attaining quite high accuracy in determining not only the velocity components, but their derivatives which determine the vertical profiles of characteristics of turbulence  $k(z)$ ,  $b(z)$ , and  $\ell(z)$  with a comparatively small number of iterations (several tens). Applying the Adams method to system (39),

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$$\left. \begin{array}{l} A'_{j+1} = A'_j - h W_1(\beta_j A_j) \\ B'_{j+1} = B'_j + h W_1(\beta_j A_j) \\ A_{j+1} = A'_j + h A'_j - h^2 W_2(\beta_j B_j) \\ B_{j+1} = B'_j + h B'_j + h^2 W_2(\beta_j A_j) \end{array} \right\} \quad (43)$$

and similar formulas for C, D, C', and D'. The following notation is used here:

$$\left. \begin{array}{l} W_1(f) = \frac{1}{60480} [198721f_j - 447288f_{j-1} + 705549f_{j-2} - \\ - 688256f_{j-3} + 407139f_{j-4} - 134472f_{j-5} + 19087f_{j-6}] \\ W_2(f) = \frac{1}{120960} [139849f_j - 243594f_{j-1} + 369399f_{j-2} - \\ - 354188f_{j-3} + 207495f_{j-4} - 68106f_{j-5} + 9625f_{j-6}] \end{array} \right\} \quad (44)$$

In order to determine the initial six values of the desired functions, it is assumed that the forces of turbulent friction noticeably overlap the effect of the Coriolis force in the surface boundary sublayer of the air, and the latter can be taken into consideration approximately in the layer  $0 < n < 6h$ .  
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$$\left. \begin{aligned} O_1(\eta) &\approx 1 + h \int_0^{\eta} \beta(t)(\eta - t) dt \\ O_2(\eta) &\approx \eta + h \int_0^{\eta} \beta(t)(\eta - t)t dt \end{aligned} \right\} \quad (45)$$

Representing the function  $\beta(n)$  by the Lagrange interpolation formula for  $j$  ( $j \leq 6$ ) of equispaced net points, the following formulas are obtained for computing the first seven ( $j = 0, 1, \dots, 6$ ) initial points:

$$\left. \begin{aligned} A'_j &= 1, \quad A_j = 0 \\ B'_j &= h^2 (\bar{\beta} \bar{b}_j), \quad B_j = h (\bar{\beta} \bar{b}_j) \\ C'_j &= h j, \quad C_j = 1 \\ D'_j &= h^3 (\bar{\beta} \bar{d}_j), \quad D_j = h^2 (\bar{\beta} \bar{d}_j) \end{aligned} \right\} \quad (46)$$

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where  $\tilde{\beta}$  is the column vector

$$\tilde{\beta} = \begin{bmatrix} \beta_0 \\ \beta_1 \\ \beta_2 \\ \beta_3 \\ \beta_4 \\ \beta_5 \\ \beta_6 \\ \beta_7 \end{bmatrix}, \quad (\text{A})$$

and  $b_j$ ,  $b_j'$ ,  $\bar{d}_j$ ,  $\bar{d}_j'$  are the  $j$ -th rows of matrices  $\bar{B}$ ,  $\bar{B}'$ ,  $\bar{D}$ , and  $\bar{D}'$ .

After computing the first seven points by formulas (46), further computations are carried out with formulas (43). After taking a sufficient number of steps and computing the values of the functions  $L$  and  $\bar{L}$  by formulas (42) for the last step  $j = N$  ( $n_h = Nh$ ), the values of the speeds  $u_0(n)$ ,  $v_0(n)$ , and their derivatives are found at points  $n_j$  ( $j = 0, 1, \dots, N$ ) by formulas (41).

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	0	0	0	0	0	0	0
1	3	6	0	0	0	0	0
2	4	0	0	0	0	0	0
3	3	0	0	0	0	0	0
39	27	27	3	0	0	0	0
40	10	40	20	0	0	0	0
56	64	16	64	0	0	0	0
45	15	15	45	0	0	0	0
1525	11 675	625	3125	625	275	0	0
1008	2016	504	1008	1008	2016	0	0
123	54	27	204	27	54	0	0
70	7	35	35	40	35	0	0
	0	0	0	0	0	0	0
1	1	0	0	0	0	0	0
2	2	0	0	0	0	0	0
3	4	1	0	0	0	0	0
3	3	3	0	0	0	0	0
3	9	9	3	0	0	0	0
8	8	8	8	0	0	0	0
14	64	8	64	14	0	0	0
45	45	15	45	45	0	0	0
05	125	125	125	125	05	0	0
288	96	144	144	96	288	0	0
41	54	27	68	27	54	41	0
	140	35	140	35	140	35	140

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(b)

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	0	0	0	0	0	0	0
	1	1	0	0	0	0	0
	12	12					
	2	16	2	0	0	0	0
	15	15	15				
	9	81	81	9	0	0	0
	40	40	40	40			
	16	1024	144	1024	16	0	0
	63	315	35	315	63		
	1375	15	95	95	15	1375	0
	4032	448	2016	2016	448	4032	
	9	972	243	72	243	972	0
	25	175	35	7	35	175	25
	0	0	0	0	0	0	0
	1	1	0	0	0	0	0
	6	3					
	4	2	0	0	0	0	0
	3	3					
	3	27	27	39	0	0	0
	20	40	10	40			
	64	16	64	56	0	0	0
	45	15	15	45			
	275	625	3125	625	11875	1525	0
	2016	1008	1008	504	2016	1008	
	54	27	204	27	54	123	0
	35	70	35	35	7	70	

(B) (Cont.)

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When  $N$  is sufficiently large, the values of  $b_N$  and  $\bar{L}_N$  vary little, which can serve as a criterion for selecting the number of iterations  $N$ .

The zero approximation for the turbulence factor and the step size  $h$  are selected. It is assumed for the first profile that  $k(z)$  is a linear function of height through the entire boundary layer:

$$k(x) = x v_0 z. \quad (47)$$

Then the following expression is obtained for  $\beta_j$ :

$$\beta_j = \frac{x^2}{m_0 R \omega} e^{\frac{x^2}{2m_0 R \omega}}. \quad (48)$$

The Rossby parameter used in formula (48) is determined by using the exact solution of the equation of motion (17) with conditions (19) and (20), where the turbulence factor is computed by formula (47):

$$G(\eta) = 1 - u_0(\eta) - l v_0(\eta) = \frac{K_0 \left( \frac{2}{\pi} \sqrt{m_0 l} e^{\frac{x^2}{2m_0 R \omega}} \right)}{K_0 \left( \frac{2}{\pi} \sqrt{m_0 l} \right)}, \quad (49)$$

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Starting with (49) and making use of  $K_0$  for small values of the argument, the parameters  $R_0$  and  $m_0$  are found to be connected by the formula

$$R_0 \approx \frac{\sqrt{\frac{x^2}{4} + \left(2c + \ln \frac{m_0}{x^2}\right)^2}}{m_0}, \quad (50)$$

where  $c = 0.5772$ , the Euler constant.

When selecting the step size  $h$ , it is assumed that the correction for the Coriolis force is very small in the layer  $0 \leq n \leq 6h$ .

On the basis of (45),

$$\delta = \int_0^{ch} \beta(t) dt, \quad (51)$$

where  $\delta$  is a small number on the order of  $10^{-2}$ — $10^{-7}$ . Actual values of  $\delta$  depend on the required accuracy. Using (51) and (48), the following transcendental equation is obtained for

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determining h:

$$e^{\frac{6\pi h}{m_0 Ro}} \left[ \frac{6\pi^2 h}{m_0 Ro} - 1 \right] = \frac{\tau_{23}}{m_0} - 1. \quad (52)$$

Thus, the order of computation is as follows:

- 1) The value of  $Ro$  is determined for the given value of the parameter  $m_0$ , then the step size  $h$  by (50);
- 2) The zero profile  $A_j$  is computed by formula (48);
- 3) The first seven values of  $A_j$ ,  $A_j'$ , etc., are computed by formulas (46);
- 4) The remaining values of  $A$  for  $j = 7, \dots, N$  are determined by formula (43);
- 5) The values of speeds  $u_0(n)$ ,  $v_0(n)$ , and their derivatives at all points  $n_j$  are computed by formulas (41), then the function  $\epsilon(n)$  by (33);
- 6) A new profile of the turbulence factor  $\theta(n)$  is determined by formula (34). If  $\theta(n)$  differs noticeably from its preceding value, all operations are repeated, starting with step 3).

When the first profile  $\theta(n)$  is computed, the formulas given

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here can be used to determine other characteristics of the boundary layer as functions of the dimensionless height  $\eta$ .

Their dependence on the real height  $z$  can be determined by the formula:

$$z = z_0 \left[ 1 + Ro \int_0^\eta \beta(\eta) d\eta \right]. \quad (53)$$

This method was tested on an M-20 computer for different values of the Rossby parameter; 5-7 approximations were required to achieve convergence of  $\beta(\eta)$  with an accuracy of several percent (see Figs. 1 and 2). Orig. art. has: 2 figures, 52 formulas and 1 table. [W.A. No. 50; ATD Report 111]

SUB CODE: 04 / SUBM DATE: none / ORIG REF: 001 / OTH REF: 001

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Next, sufficiently  
stable, which can occur  
under certain conditions,  
the approximation  
size of the selected  
height is linear  
entry layer.

the following way:  
First, by parameter  
of the transition  
conditions ( $\log \frac{z}{z_0}$ ) are  
calculated, the

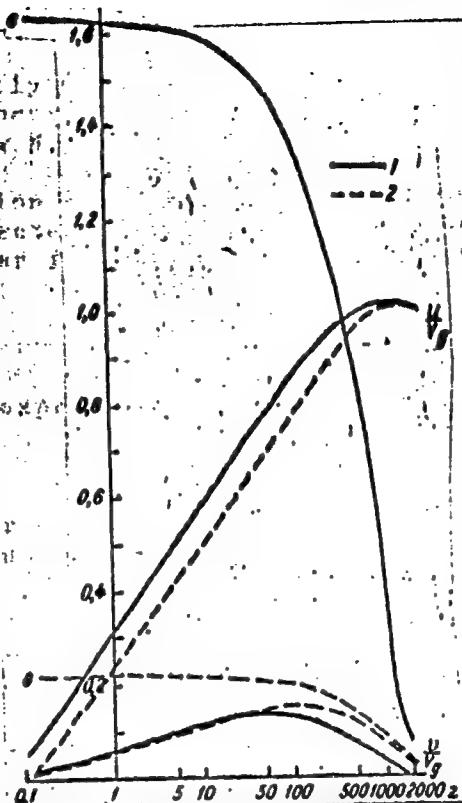


Fig. 1. Vertical profiles  
of the wind velocity compo-  
nents  $u/v_g$ ,  $v/v_g$ , and  
turbulent energy  $b(z)$ .

1 -  $\log \text{Ro} = 6.1$ ; 2 -  $\log$   
 $\text{Ro} = 5.6$

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where  $\rho \approx 10^{-3} \text{ kg/cm}^3$ , the  
mean velocity  $U_0 = 10 \text{ cm/sec}$ ,  
the rotation  $\Omega = 10^{-3} \text{ sec}^{-1}$ ,  
and  $R_0 = 10 \text{ cm}$ .

On the basis of these

values we find

that  $\log R_o = 6.1$  and

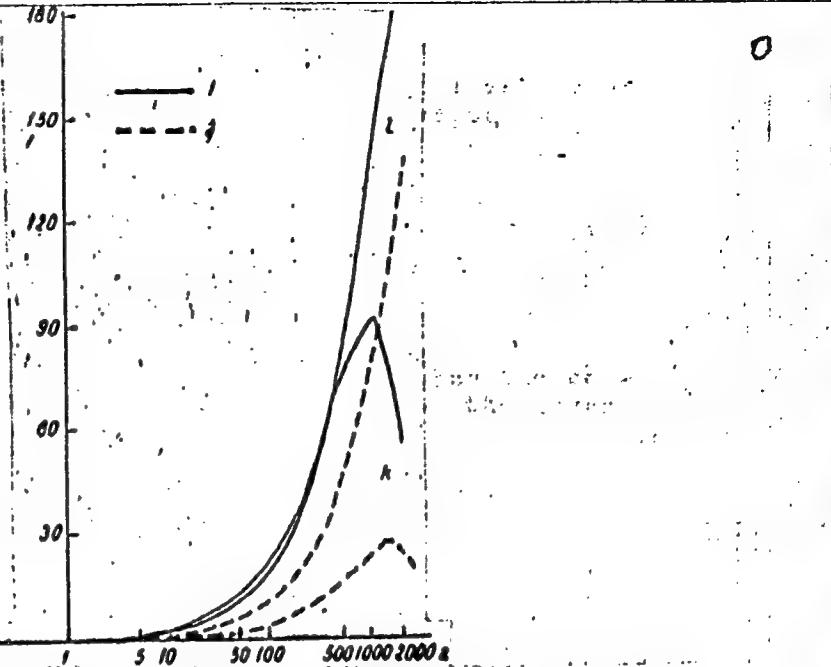
$\log \Omega = 5.6$ .

Fig. 2. Vertical profiles of the turbulence factor

$k(z)$  and the scale of turbulence  $\ell(z)$

$1 - \log R_o = 6.1; 2 - \log R_o = 5.6.$

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AP6032071

SOURCE CODE: UR/0362/66/002/009/0920/0927

AUTHOR: Vager, B. G.

ORG: Main Geophysical Observatory (Glavnaya geofizicheskaya observatoriya)

TITLE: Effect of turbulent diffusion in a semi-empirical model of the lower atmosphere

SOURCE: AN SSSR Izvestiya. Fizika atmosfery i okeana, v. 2, no. 9, 1966, 920-927

TOPIC TAGS: lower atmosphere, atmospheric turbulence, atmospheric model, atmospheric diffusion, energy diffusion, turbulent energy

ABSTRACT: The author investigated a semi-empirical model of turbulent transfer in the lower atmosphere proposed by S. S. Zilitinkevich and D. D. Laykhtman. A numerical solution of the problem is given, taking into consideration the effect of turbulent energy diffusion. It is shown that the effect plays an important role in unstable atmospheric stratification. The author thanks, S. S. Zilitinkevich for his help and valuable advice. Orig. art. has: 4 figures, 23 formulas and 1 table. [Based on author's abstract] SUB CODE: 04/SUBM DATE:22Apr66/ORIG REF: 009/

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OTH REF: 001/

UDC: 551.551.8

25(5)

SOV/28-59-2-14/36

AUTHORS: Rabinovich, P.M., Khrisanfov, G.A.,( Moscow) Tager, I.A.  
(Moscow) and Shitikov, A.M.,(Leningrad); Engineers.

TITLE: On Revising Standards of the "Layout System" (K peresmotru  
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PERIODICAL: Standartizatsiya, 1959, Nr 2, pp 43-47 (USSR)

ABSTRACT: This article contains suggestions by four authors for lay-  
out standards now being revised. Different modifications  
to the preparation of working drawings, their registration  
and their storage, are proposed. There are 2 tables.

ASSOCIATION: TsNIITMASH; VNII.

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Mar/Apr 53

"Stolbur Wilting of Potatoes", R.M. Vager, Moscow Sta, All-Union Inst of Plant Protection

Mikrobiologiya, Vol 22, No 2, pp 198-202.

Stolbur (caused by *H. obsoletus*) is entirely distinct from fusariosis and independent of secondary infection with the latter disease. In combatting potato wilting, one must pay primary attention to stolbur. In the experiments described, infection with *H. obsoletus* was carried out by using insects as transmitters or by applying the method of grafting.

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E.

Abs Jour : Ref Zhur - Biol., No 19, 1958, 85780

Author : Vager, R.M.

Inst :

Title : The Production of Specific Serum Against the Mosaic Virus  
of Winter Wheat.

Orig Pub : Dokl. VASKhNIL, 1957, No 12, 20-21

Abstract : Serum was prepared against the virus. The disease is transmitted by cicadas and is classified with the plant diseases known as the yellows. The total amount of injected antigen comprised 69.6 mg in doses of 2.5 to 6 ml of liquid. The serum produced a reaction only with the juice of sick plants in a dilution of 1:8.

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(BASAL METABOLISM pharmacol) (IODINE radioactive)

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Budapest, Medicai Akadémiai Kiadó, 1960, vol. 6, pp. 66-69.

Abstract: [Authors' English summary modified] The bactericidal effect of a weak aqueous solution of chavicol-methylxylylum 1: cream ( $C_6H_9CH_2O_2$ ) - phenoxymethane (i.e. K. I.) was studied experimentally, with special emphasis on the bactericidal effect of the solution when prepared by the 0.2 percent solution of phenoxymethane (i.e. 0.1% of Phenox) even in a 10-fold dilution. In laboratory researches no effect was observed and the solution was harmless to the intestinal mucosa in contact with it. When larger quantities were introduced, however, (i.e. for rinsing the abdominal cavity) big poisoning developed. Following i.v. administration, the experimental animal died; when the animal was killed it was found that the solution had caused a large amount of edema in the liver, grossly purulent infection of the liver, necrosis of the liver, and edema and hemorrhage and the fibrin was found in the liver, lungs before reposition, will be mentioned.

1/1

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Budapest, Hungary

Re: East European Accession, Vol. 5, No. 5, May 1956

HUNGARY

FELKAI, Ferenc, Dr., and VAGHY, Tamas, Dr., Assistant Professors, Chair and Clinic for Surgery and Ophthalmology at the University for Veterinary Sciences (Allatorvostudomanyi Egyetem Sebeszeti es Szemészeti Tanszeke es Klinikaja) [location not given] (Head: KOVACS, Andras B., Dr., Professor, Candidate of Veterinary Sciences).

"Some Instances of the Application of Trypsin Wound Powder"

Budapest, Mazsar Allatorvosok Lapja, Vol 21, No 5, May 1966, pp 231-235.

Abstract: Twenty-six cases illustrating the uses of Trypsin wound powder, containing 0.25 g. trypsin, local anesthetic, stabilizer, and glycocoll in 10-g. shaker-type jars, were described. The experiences were generally favorable, confirming reports in the literature regarding the performance of trypsin-containing wound powders. It was especially effective in the treatment of abscesses and fistulas. No allergenic reactions were noted. The powder should be administered every 6-8 hours for about 2-3 days. It does not serve as a substitute for medical or surgical procedures, but it supplements them. 29 references, including 8 Hungarian, 4 German, and 17 Western.

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VACHY, Tamas, Dr, MESZAROS, Janos, Dr; Veterinary Medical University, Department of Surgery and Ophthalmology, and Clinic (chairman: B., KOVACS, Andras, Dr, professor, cand. of vet. sci.) and Department of Pathological Anatomy (chairman: KARDEVAN, Andor, Dr, docent) (Allatorvostudomanyi Egyetem, Sebeszeti es Szemeszeti Tanszek es Klinikai es Korbonctani Tanszek).

"Intermittent Lameness in a Horse Caused by Thrombosis."

Budapest, Magyar Allatorvosok Lapja, Vol 21, No 10, Oct 66, pages 461-463.

Abstract: [Authors' English summary modified] On the basis of the clinical symptoms and of the results of rectal and other complementary examinations (pCO<sub>2</sub>determination of the venous blood), the presence of a thrombi was diagnosed in the quadrifurcation of the aorta as well as in the pelvic and iliac arteries. The animal received no treatment because of the poor prognosis. In the course of dissection of the animal, large white and mixed thrombi were found in the above mentioned vessels which filled up the lumen to about three-fourth of its volume. It is suggested by the authors that the pathological process originated most probably from an unknown injury to the intima. The post mortem examination revealed that the development of the thrombosis was started in the right iliac artery. Chondrous islands originating from metaplasia were also found at several places of the organized areas. 8 Hungarian, 15 Western references.

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VAGI, Ferenc, dr., kandidatus, tanszekvezeto docens

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"Economy in Road Planning; Remarks on Istvan Czuni's Study", P. 104,  
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Hungary)

SO: Monthly List of East European Accessions, (EEAL), LC, Vol. 3, No. 12,  
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Melyepitestudomanyi Szemle. - Vol. 4, no. 12, Dec. 1954.

Rounding off longitudinal sections and building superelevated crossings on public roads.  
p. 649.

SO: Monthly list of East European Accessions, (EEAL), LC, Vol. 4, No. 9, Sept. 1955  
Uncl.

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Temporary roads to construction sites. p. 543.  
Vol 5, no. 12, Dec. 1955. MÉLYÉPÍTÉSTUDOMÁNYI SZEMLE. Budapest, Hungary.

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Budapest, Hungary

So: East European Acquisitions, Vol. 5, no. 5, May 1956

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Impressions from travel in the German Democratic Republic, p. 369,  
MELYEPITESTUDOMANYI SZEMLE (Kozlekedesi Kiado) Budapest, Vol. 6,  
No. 7/8, July/Aug. 1956

SOURCE: East European Accessions List (EEAL) Library of Congress,  
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VAGI, Istvan

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